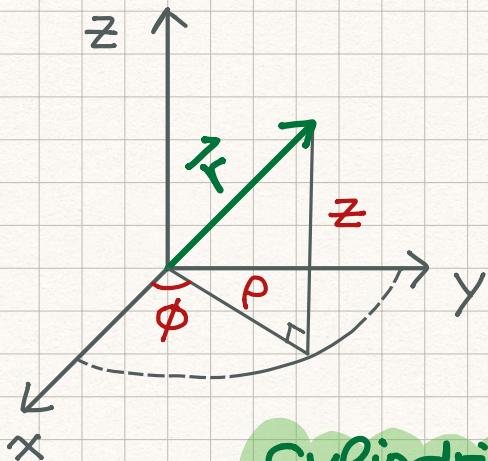


# Curvilinear Coordinates



Cylindrical  
coordinates

The "del"  $\vec{\nabla}$  operator takes rather different forms in the curvilinear coordinates  $\vec{r}$

$$\begin{aligned}x &= \rho \cos\phi \\y &= \rho \sin\phi \\z &= z\end{aligned}$$

$$\begin{aligned}\text{The position vector } \vec{r} &= x\hat{i} + y\hat{j} + z\hat{k} \\&= \rho \cos\phi \hat{i} + \rho \sin\phi \hat{j} + z\hat{k}\end{aligned}$$

$$\vec{a}_\rho = \frac{\partial \vec{r}}{\partial \rho} = \cos\phi \hat{i} + \sin\phi \hat{j}$$

$$\vec{a}_\phi = \frac{\partial \vec{r}}{\partial \phi} = -\rho \sin\phi \hat{i} + \rho \cos\phi \hat{j}$$

$$\vec{a}_z = \frac{\partial \vec{r}}{\partial z} = \hat{k}$$

$\vec{a}_\rho, \vec{a}_\phi, \vec{a}_z$  are the base vectors for  $(\rho, \phi, z)$

normalize to unit vectors :

$$\hat{e}_\rho = \vec{a}_\rho = \cos\phi \hat{i} + \sin\phi \hat{j}$$

$$\hat{e}_\phi = \frac{1}{\rho} \vec{a}_\phi = -\sin\phi \hat{i} + \cos\phi \hat{j}$$

$$\hat{e}_z = \vec{a}_z = \hat{k}$$

$\{\hat{e}_\rho, \hat{e}_\phi, \hat{e}_z\}$  orthonormal basis

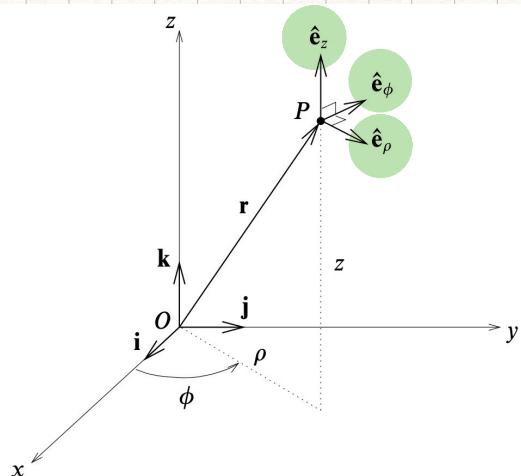


Figure 10.7 Cylindrical polar coordinates  $\rho, \phi, z$ .

The infinitesimal displacement  $d\vec{r}$  is

$$\begin{aligned} d\vec{r} &= \frac{\partial \vec{r}}{\partial p} dp + \frac{\partial \vec{r}}{\partial \phi} d\phi + \frac{\partial \vec{r}}{\partial z} dz \\ &= dp \hat{\vec{a}}_p + d\phi \hat{\vec{a}}_\phi + dz \hat{\vec{a}}_z \\ &= dp \hat{\vec{e}}_p + \rho d\phi \hat{\vec{e}}_\phi + dz \hat{\vec{e}}_z \end{aligned}$$

The element of arc length  $ds$  is

$$(ds)^2 = d\vec{r} \cdot d\vec{r} = (dp)^2 + \rho^2 (d\phi)^2 + (dz)^2$$

**NOTE.**

In general curvilinear coordinates  $q_i$ ,

$$\hat{\vec{a}}_i = \frac{\partial \vec{r}}{\partial q_i} \quad \leftarrow \text{not necessarily orthogonal.}$$

$$(ds)^2 = d\vec{r} \cdot d\vec{r} = \sum_i \frac{\partial \vec{r}}{\partial q_i} dq_i \cdot \sum_j \frac{\partial \vec{r}}{\partial q_j} dq_j$$

$$= \sum_{ij} \frac{\partial \vec{r}}{\partial q_i} \cdot \frac{\partial \vec{r}}{\partial q_j} dq_i dq_j$$

$\hookrightarrow$  metric tensor  $g_{ij} = \hat{\vec{a}}_i \cdot \hat{\vec{a}}_j$

$$(ds)^2 = \sum_{ij} g_{ij} dq_i dq_j$$

describes the curvature  
of the space

For the cylindrical,

$$g_{ij} = 0, \quad i \neq j$$

$$g_{ij} = h_i^2 \delta_{ij}$$

$$h_p = 1, \quad h_\phi = \rho, \quad h_z = 1$$

# Spherical Coordinates

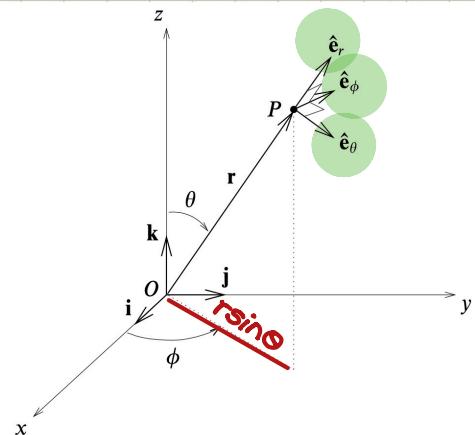


Figure 10.9 Spherical polar coordinates  $r, \theta, \phi$ .

$$\begin{aligned}x &= r \sin \theta \cos \phi \\y &= r \sin \theta \sin \phi \\z &= r \cos \theta\end{aligned}$$

$$\begin{aligned}\vec{r} &= x \hat{i} + y \hat{j} + z \hat{k} \\&= r \sin \theta \cos \phi \hat{i} \\&\quad + r \sin \theta \sin \phi \hat{j} \\&\quad + r \cos \theta \hat{k}\end{aligned}$$

The base vectors  $\vec{a}_i$  are

$$\vec{a}_r = \frac{\partial \vec{r}}{\partial r} = \sin \theta \cos \phi \hat{i} + \sin \theta \sin \phi \hat{j} + \cos \theta \hat{k}$$

$$\vec{a}_\theta = \frac{\partial \vec{r}}{\partial \theta} = r \cos \theta \cos \phi \hat{i} + r \cos \theta \sin \phi \hat{j} - r \sin \theta \hat{k}$$

$$\vec{a}_\phi = \frac{\partial \vec{r}}{\partial \phi} = -r \sin \theta \sin \phi \hat{i} + r \sin \theta \cos \phi \hat{j}$$

$$|\vec{a}_r| = 1, \quad |\vec{a}_\theta| = r, \quad |\vec{a}_\phi| = r \sin \theta$$



normalize to unit vectors,

$$\hat{e}_r = \vec{a}_r = \sin \theta \cos \phi \hat{i} + \sin \theta \sin \phi \hat{j} + \cos \theta \hat{k}$$

$$\hat{e}_\theta = \frac{1}{r} \vec{a}_\theta = \cos \theta \cos \phi \hat{i} + \cos \theta \sin \phi \hat{j} - \sin \theta \hat{k}$$

$$\hat{e}_\phi = \frac{1}{r \sin \theta} \vec{a}_\phi = -\sin \phi \hat{i} + \cos \phi \hat{j}$$

$\{\hat{e}_r, \hat{e}_\theta, \hat{e}_\phi\}$  orthonormal basis

The element of arc length  $dS$  is

$$\begin{aligned}
 (dS)^2 &= \sum_{ij} g_{ij} dq_i dq_j - g_{ij} = \vec{a}_i \cdot \vec{a}_j \\
 &= \sum_i h_i^2 (dq_i)^2 - h_i = |\vec{a}_i| \text{ scale factor} \\
 &= (dr)^2 + r^2(d\theta)^2 + r^2 \sin^2\theta (d\phi)^2
 \end{aligned}$$

scale factors

$$h_r = 1, h_\theta = r, h_\phi = r \sin\theta$$

singular @  $r=0$

singular  
@  $r=0$   
@  $\theta=0, \pi$

## volume element

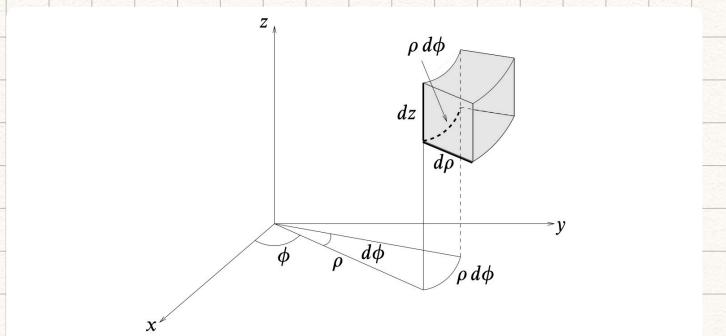


Figure 10.8 The element of volume in cylindrical polar coordinates is given by  $\rho d\rho d\phi dz$ .

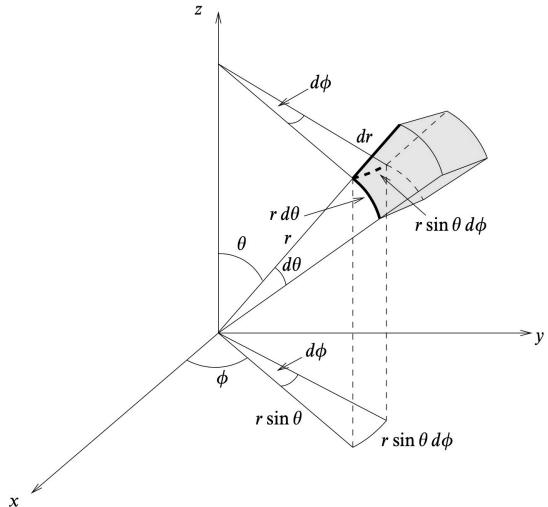
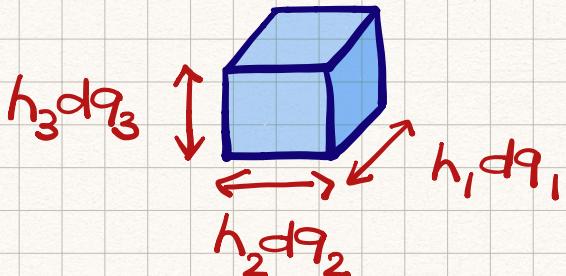


Figure 10.10 The element of volume in spherical polar coordinates is given by  $r^2 \sin\theta dr d\theta d\phi$ .



Mr. Cube

$$d\tau = h_1 h_2 h_3 dq_1 dq_2 dq_3$$

# Gradient

By definition,  $d\Phi = \vec{\nabla} \Phi \cdot d\vec{r}$

By partial differentiation,  $d\Phi = \sum_{i=1}^3 \frac{\partial \Phi}{\partial q_i} dq_i$

Set  $dq_2 = dq_3 = 0$ , — holding  $q_2, q_3$  constant.

$$d\vec{r} = \frac{\partial \vec{r}}{\partial q_1} dq_1 = \hat{e}_1 h_1 dq_1$$

$$\text{Thus, } d\Phi = \vec{\nabla} \Phi \cdot \hat{e}_1 h_1 dq_1 = \frac{\partial \Phi}{\partial q_1} dq_1$$



$$\vec{\nabla} \Phi \cdot \hat{e}_1 = \frac{1}{h_1} \frac{\partial \Phi}{\partial q_1}$$

Similarly, one can work out the other comp.

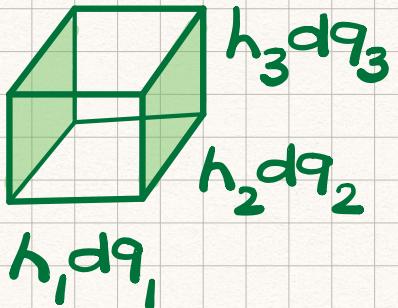
$$\vec{\nabla} \Phi = \frac{1}{h_1} \frac{\partial \Phi}{\partial q_1} \hat{e}_1 + \frac{1}{h_2} \frac{\partial \Phi}{\partial q_2} \hat{e}_2 + \frac{1}{h_3} \frac{\partial \Phi}{\partial q_3} \hat{e}_3$$

OR, in the operator form ~

$$\vec{\nabla} = \hat{e}_1 \frac{1}{h_1} \frac{\partial}{\partial q_1} + \hat{e}_2 \frac{1}{h_2} \frac{\partial}{\partial q_2} + \hat{e}_3 \frac{1}{h_3} \frac{\partial}{\partial q_3}$$

# Divergence

Mr. Cube



$$\int_{\partial V} \vec{J} \cdot d\vec{a} = \int_V \vec{\nabla} \cdot \vec{J} d\tau$$

Flux along  $q_1$  direction :

$$\begin{aligned}\Phi_1 &= J_1 h_2 h_3 (q_1 + dq_1) dq_2 dq_3 \\ &\quad - J_1 h_2 h_3 (q_1) dq_2 dq_3 \\ &= \frac{\partial}{\partial q_1} (h_2 h_3 J_1) dq_1 dq_2 dq_3\end{aligned}$$

$$\int_{\partial V} \vec{J} \cdot d\vec{a} = \Phi_1 + \Phi_2 + \Phi_3$$

$$= \left[ \frac{\partial}{\partial q_1} (h_2 h_3 J_1) + \frac{\partial}{\partial q_2} (h_3 h_1 J_2) + \frac{\partial}{\partial q_3} (h_1 h_2 J_3) \right] dq_1 dq_2 dq_3$$

The volume integral is simple ~

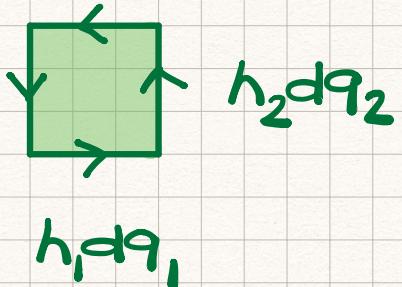
$$\int_V \vec{\nabla} \cdot \vec{J} d\tau = \vec{\nabla} \cdot \vec{J} \cdot h_1 h_2 h_3 dq_1 dq_2 dq_3$$

By comparison, it's easy to see that

$$\begin{aligned}\vec{\nabla} \cdot \vec{J} &= \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial q_1} (h_2 h_3 J_1) \right. \\ &\quad + \frac{\partial}{\partial q_2} (h_1 h_3 J_2) \\ &\quad \left. + \frac{\partial}{\partial q_3} (h_1 h_2 J_3) \right]\end{aligned}$$

# Curl

Ms Loop



$$\oint_{\partial S} \vec{V} \cdot d\vec{r} = \int_S \vec{\nabla} \times \vec{V} \cdot d\vec{S}$$

Consider the circulation in the  $q_1-q_2$  plane. It contains all contributions from the 4 segments :

$$\begin{aligned} \oint_{\partial S} \vec{V} \cdot d\vec{r} &= h_2 V_2 (q_1 + dq_1) dq_2 - h_2 V_2 (q_1) dq_2 \\ &\quad - h_1 V_1 (q_2 + dq_2) dq_1 + h_1 V_1 (q_2) dq_1 \\ &= \left[ \frac{\partial}{\partial q_1} (h_2 V_2) - \frac{\partial}{\partial q_2} (h_1 V_1) \right] dq_1 dq_2 \end{aligned}$$

The surface integral can be expressed as

$$\int_S \vec{\nabla} \times \vec{V} \cdot d\vec{S} = (\vec{\nabla} \times \vec{V}) \cdot \hat{e}_3 h_1 h_2 dq_1 dq_2$$

By comparison, the curl  $(\vec{\nabla} \times \vec{V})_3$  is

$$(\vec{\nabla} \times \vec{V}) \cdot \hat{e}_3 = \frac{1}{h_1 h_2} \left[ \frac{\partial}{\partial q_1} (h_2 V_2) - \frac{\partial}{\partial q_2} (h_1 V_1) \right]$$

working  
all comp.



$$\vec{\nabla} \times \vec{V} = \frac{1}{h_1 h_2 h_3}$$

$h_1 \hat{e}_1$	$h_2 \hat{e}_2$	$h_3 \hat{e}_3$
$\frac{\partial}{\partial q_1}$	$\frac{\partial}{\partial q_2}$	$\frac{\partial}{\partial q_3}$
$h_1 V_1$	$h_2 V_2$	$h_3 V_3$

# Summary

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$$\nabla \Phi = \frac{1}{h_1} \frac{\partial \Phi}{\partial u_1} \hat{\mathbf{e}}_1 + \frac{1}{h_2} \frac{\partial \Phi}{\partial u_2} \hat{\mathbf{e}}_2 + \frac{1}{h_3} \frac{\partial \Phi}{\partial u_3} \hat{\mathbf{e}}_3$$

$$\nabla \cdot \mathbf{a} = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} (h_2 h_3 a_1) + \frac{\partial}{\partial u_2} (h_3 h_1 a_2) + \frac{\partial}{\partial u_3} (h_1 h_2 a_3) \right]$$

$$\nabla \times \mathbf{a} = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \hat{\mathbf{e}}_1 & h_2 \hat{\mathbf{e}}_2 & h_3 \hat{\mathbf{e}}_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ h_1 a_1 & h_2 a_2 & h_3 a_3 \end{vmatrix}$$

$$\nabla^2 \Phi = \frac{1}{h_1 h_2 h_3} \left[ \frac{\partial}{\partial u_1} \left( \frac{h_2 h_3}{h_1} \frac{\partial \Phi}{\partial u_1} \right) + \frac{\partial}{\partial u_2} \left( \frac{h_3 h_1}{h_2} \frac{\partial \Phi}{\partial u_2} \right) + \frac{\partial}{\partial u_3} \left( \frac{h_1 h_2}{h_3} \frac{\partial \Phi}{\partial u_3} \right) \right]$$


---

Table 10.4 Vector operators in orthogonal curvilinear coordinates  $u_1, u_2, u_3$ .  
 $\Phi$  is a scalar field and  $\mathbf{a}$  is a vector field.

cylindrical

$$h_\rho = 1, \quad h_\phi = \rho, \quad h_z = 1$$

---


$$\nabla \Phi = \frac{\partial \Phi}{\partial \rho} \hat{\mathbf{e}}_\rho + \frac{1}{\rho} \frac{\partial \Phi}{\partial \phi} \hat{\mathbf{e}}_\phi + \frac{\partial \Phi}{\partial z} \hat{\mathbf{e}}_z$$

$$\nabla \cdot \mathbf{a} = \frac{1}{\rho} \frac{\partial}{\partial \rho} (\rho a_\rho) + \frac{1}{\rho} \frac{\partial a_\phi}{\partial \phi} + \frac{\partial a_z}{\partial z}$$

$$\nabla \times \mathbf{a} = \frac{1}{\rho} \begin{vmatrix} \hat{\mathbf{e}}_\rho & \rho \hat{\mathbf{e}}_\phi & \hat{\mathbf{e}}_z \\ \frac{\partial}{\partial \rho} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ a_\rho & \rho a_\phi & a_z \end{vmatrix}$$

$$\nabla^2 \Phi = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left( \rho \frac{\partial \Phi}{\partial \rho} \right) + \frac{1}{\rho^2} \frac{\partial^2 \Phi}{\partial \phi^2} + \frac{\partial^2 \Phi}{\partial z^2}$$


---

Table 10.2 Vector operators in cylindrical polar coordinates;  $\Phi$  is a scalar field and  $\mathbf{a}$  is a vector field.

spherical

$$\lambda_r = 1, \quad \lambda_\theta = r, \quad \lambda_\phi = r \sin \theta$$

$$\nabla \Phi = \frac{\partial \Phi}{\partial r} \hat{\mathbf{e}}_r + \frac{1}{r} \frac{\partial \Phi}{\partial \theta} \hat{\mathbf{e}}_\theta + \frac{1}{r \sin \theta} \frac{\partial \Phi}{\partial \phi} \hat{\mathbf{e}}_\phi$$

$$\nabla \cdot \mathbf{a} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 a_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta a_\theta) + \frac{1}{r \sin \theta} \frac{\partial a_\phi}{\partial \phi}$$

$$\nabla \times \mathbf{a} = \frac{1}{r^2 \sin \theta} \begin{vmatrix} \hat{\mathbf{e}}_r & r \hat{\mathbf{e}}_\theta & r \sin \theta \hat{\mathbf{e}}_\phi \\ \frac{\partial}{\partial r} & \frac{\partial}{\partial \theta} & \frac{\partial}{\partial \phi} \\ a_r & r a_\theta & r \sin \theta a_\phi \end{vmatrix}$$

$$\nabla^2 \Phi = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \Phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \Phi}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2}$$

Table 10.3 Vector operators in spherical polar coordinates;  $\Phi$  is a scalar field and  $\mathbf{a}$  is a vector field.