



國立清華大學

Electromagnetism

Introduction to Electrodynamics 4th David J. Griffiths

Chap.10 Potentials and Fields

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Exercise List

4, 11, 15, 16, 27, 28

Problem 10.4 Suppose $V = 0$ and $\mathbf{A} = A_0 \sin(kx - \omega t) \hat{\mathbf{y}}$, where A_0 , ω , and k are constants. Find \mathbf{E} and \mathbf{B} , and check that they satisfy Maxwell's equations in vacuum. What condition must you impose on ω and k ?

$$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t} = -0 + \omega A_0 \cos(kx - \omega t) \hat{\mathbf{y}}$$

$$\mathbf{B} = \nabla \times \mathbf{A} = k A_0 \cos(kx - \omega t) \hat{\mathbf{z}}$$

In vacuum $\rho = 0$ & $\mathbf{J} = 0$

$$\nabla \cdot \mathbf{E} = 0 \Rightarrow \nabla \cdot \mathbf{E} = \partial_y \omega A_0 \cos(kx - \omega t) = 0$$

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \Rightarrow \begin{cases} \nabla \times \mathbf{E} = \partial_x \omega A_0 \cos(kx - \omega t) \hat{\mathbf{z}} = -k \omega A_0 \sin(kx - \omega t) \hat{\mathbf{z}} \\ \frac{\partial \mathbf{B}}{\partial t} = k \omega A_0 \sin(kx - \omega t) \hat{\mathbf{z}} \end{cases}$$

$$\nabla \cdot \mathbf{B} = 0 \Rightarrow \nabla \cdot \mathbf{B} = \partial_z k A_0 \cos(kx - \omega t) = 0$$

$$\nabla \times \mathbf{B} = \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t} \Rightarrow \begin{cases} \nabla \times \mathbf{B} = -\partial_x k A_0 \cos(kx - \omega t) \hat{\mathbf{y}} = k^2 A_0 \sin(kx - \omega t) \hat{\mathbf{y}} \\ \frac{\partial \mathbf{E}}{\partial t} = \omega^2 A_0 \sin(kx - \omega t) \hat{\mathbf{y}} \end{cases}$$

$$k^2 = \mu_0 \epsilon_0 \omega^2 \Rightarrow \omega = \frac{k}{\sqrt{\mu_0 \epsilon_0}} = ck$$

Problem 10.11

(a) Suppose the wire in Ex. 10.2 carries a linearly increasing current

$$I(t) = kt,$$

for $t > 0$. Find the electric and magnetic fields generated.

It's electrically neutral $\Rightarrow V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}', t_t)}{r} d\tau = 0$

For $t < s/c$, $\mathbf{A} = 0$ For $t > s/c$:

$$\begin{aligned} \mathbf{A}(\mathbf{r}, t) &= \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}', t_r)}{r} d\tau' = \hat{\mathbf{z}} \frac{\mu_0}{4\pi} \int_{-\sqrt{(ct)^2 - s^2}}^{\sqrt{(ct)^2 - s^2}} \frac{kt_r}{\sqrt{s^2 + z'^2}} dz' = \hat{\mathbf{z}} \left(\frac{\mu_0 k}{4\pi} \right) 2 \int_0^{\sqrt{(ct)^2 - s^2}} \frac{t - \sqrt{s^2 + z'^2}/c}{\sqrt{s^2 + z'^2}} dz' \\ &= \hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left[t \int_0^{\sqrt{(ct)^2 - s^2}} \frac{1}{\sqrt{s^2 + z'^2}} dz' - \frac{1}{c} \int_0^{\sqrt{(ct)^2 - s^2}} dz' \right] \\ &= \hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left[t \ln \left(\sqrt{s^2 + z'^2} + z' \right) \Big|_0^{\sqrt{(ct)^2 - s^2}} - \frac{\sqrt{(ct)^2 - s^2}}{c} \right] \\ &= \hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left(t \ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} - \frac{\sqrt{(ct)^2 - s^2}}{c} \right) \end{aligned}$$

Problem 10.11

(a) Suppose the wire in Ex. 10.2 carries a linearly increasing current

$$I(t) = kt,$$

for $t > 0$. Find the electric and magnetic fields generated.

$$\begin{aligned}\mathbf{E} &= -\nabla V - \frac{\partial \mathbf{A}}{\partial t} = -\hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left[\ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} + \frac{st}{ct + \sqrt{(ct)^2 - s^2}} \frac{1}{s} \left(c + \frac{c^2 t}{\sqrt{(ct)^2 - s^2}} \right) - \frac{c^2 t}{c\sqrt{(ct)^2 - s^2}} \right] \\ &= -\hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left[\ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} + \frac{ct}{ct + \sqrt{(ct)^2 - s^2}} \left(\frac{\sqrt{(ct)^2 - s^2} + ct}{\sqrt{(ct)^2 - s^2}} \right) - \frac{ct}{\sqrt{(ct)^2 - s^2}} \right] \\ &= -\hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} \quad \text{or} \quad = 0 (t < s/c)\end{aligned}$$

$$\mathbf{A} = \hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left(t \ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} - \frac{\sqrt{(ct)^2 - s^2}}{c} \right)$$

Problem 10.11

(a) Suppose the wire in Ex. 10.2 carries a linearly increasing current

$$I(t) = kt,$$

for $t > 0$. Find the electric and magnetic fields generated.

$$\mathbf{A} = \hat{\mathbf{z}} \frac{\mu_0 k}{2\pi} \left(t \ln \frac{ct + \sqrt{(ct)^2 - s^2}}{s} - \frac{\sqrt{(ct)^2 - s^2}}{c} \right)$$

$$\begin{aligned} \mathbf{B} &= \nabla \times \mathbf{A} = -\partial_s A_z \hat{\phi} = -\hat{\phi} \frac{\mu_0 k}{2\pi} \left[t \frac{s}{ct + \sqrt{(ct)^2 - s^2}} \left(-\frac{ct + \sqrt{(ct)^2 - s^2}}{s^2} - \frac{s}{s\sqrt{(ct)^2 - s^2}} \right) + \frac{s}{c\sqrt{(ct)^2 - s^2}} \right] \\ &= -\hat{\phi} \frac{\mu_0 k}{2\pi} \left[t \frac{s}{ct + \sqrt{(ct)^2 - s^2}} \left(-\frac{ct\sqrt{(ct)^2 - s^2} + (ct)^2 - s^2 + s^2}{s^2\sqrt{(ct)^2 - s^2}} \right) + \frac{s}{c\sqrt{(ct)^2 - s^2}} \right] \\ &= -\hat{\phi} \frac{\mu_0 k}{2\pi} \left[t \frac{-ct}{s\sqrt{(ct)^2 - s^2}} + \frac{s}{c\sqrt{(ct)^2 - s^2}} \right] \\ &= -\hat{\phi} \frac{\mu_0 k}{2\pi} \frac{-(ct)^2 + s^2}{sc\sqrt{(ct)^2 - s^2}} = \frac{\mu_0 k}{2\pi sc} \sqrt{(ct)^2 - s^2} \hat{\phi} \quad \text{or} \quad = 0 (t < s/c) \end{aligned}$$

Problem 10.11

(b) Do the same for the case of a sudden burst of current:

$$I(t) = q_0 \delta(t).$$

$$\begin{aligned} \mathbf{A}(\mathbf{r}, t) &= \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}', t_r)}{r} d\tau' = \hat{\mathbf{z}} \frac{\mu_0}{4\pi} \int_{-\infty}^{\infty} \frac{q_0 \delta(t_r)}{\sqrt{s^2 + z'^2}} dz' = \hat{\mathbf{z}} \frac{\mu_0 q_0}{4\pi} 2 \int_0^{\infty} \frac{\delta\left(t - \sqrt{s^2 + z'^2}/c\right)}{\sqrt{s^2 + z'^2}} dz' \\ &= \hat{\mathbf{z}} \frac{\mu_0 q_0}{2\pi} \int_s^{\infty} \frac{\delta\left(t - \sqrt{s^2 + z'^2}/c\right)}{\sqrt{s^2 + z'^2}} \left(\frac{c\sqrt{s^2 + z'^2}}{z'} d\sqrt{s^2 + z'^2}/c \right) \\ &= \frac{\mu_0 q_0}{2\pi} \frac{c}{\sqrt{(ct)^2 - s^2}} \hat{\mathbf{z}} \end{aligned}$$

$$\mathbf{B} = \nabla \times \mathbf{A} = -\partial_s A_z \hat{\phi} = -\frac{\mu_0 q_0}{2\pi} \frac{cs}{\left[(ct)^2 - s^2\right]^{3/2}} \hat{\phi} \quad \text{or} \quad = 0(t < s/c)$$

$$\mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t} = \frac{\mu_0 q_0}{2\pi} \frac{c^3 t}{\left[(ct)^2 - s^2\right]^{3/2}} \hat{\mathbf{z}} \quad \text{or} \quad = 0(t < s/c)$$

Problem 10.15 A particle of charge q moves in a circle of radius a at constant angular velocity ω . (Assume that the circle lies in the xy plane, centered at the origin, and at time $t = 0$ the charge is at $(a, 0)$, on the positive x axis.) Find the Liénard-Wiechert potentials for points on the z axis.

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{(rc - \mathbf{r} \cdot \mathbf{v})}, \quad (10.46)$$

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \frac{qc\mathbf{v}}{(rc - \mathbf{r} \cdot \mathbf{v})} = \frac{\mathbf{v}}{c^2} V(\mathbf{r}, t). \quad (10.47)$$

$$\mathbf{w}(t) = a[\cos(\omega t)\hat{\mathbf{x}} + \sin(\omega t)\hat{\mathbf{y}}] \Rightarrow \mathbf{v}(t) = \omega a[-\sin(\omega t)\hat{\mathbf{x}} + \cos(\omega t)\hat{\mathbf{y}}]$$

$$\mathbf{r} = z\hat{\mathbf{z}} - \mathbf{w}(t_r) = z\hat{\mathbf{z}} - a[\cos(\omega t_r)\hat{\mathbf{x}} + \sin(\omega t_r)\hat{\mathbf{y}}], \quad r = \sqrt{z^2 + a^2} = c(t - t_r) \Rightarrow t_r = t - \frac{\sqrt{z^2 + a^2}}{c}$$

$$\mathbf{r} \cdot \mathbf{v}(t_r) = \omega a^2 \sin(\omega t_r) \cos(\omega t_r) (1 - 1) = 0$$

$$\left\{ \begin{array}{l} V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{(rc - \mathbf{r} \cdot \mathbf{v})} = \frac{1}{4\pi\epsilon_0} \frac{q}{\sqrt{z^2 + a^2}} \\ \mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \frac{qc\mathbf{v}}{(rc - \mathbf{r} \cdot \mathbf{v})} = \frac{\mu_0}{4\pi} \frac{q\omega a}{\sqrt{z^2 + a^2}} [-\sin(\omega t_r)\hat{\mathbf{x}} + \cos(\omega t_r)\hat{\mathbf{y}}] = \frac{1}{4\pi\epsilon_0} \frac{q\omega a}{c^2 \sqrt{z^2 + a^2}} [-\sin(\omega t_r)\hat{\mathbf{x}} + \cos(\omega t_r)\hat{\mathbf{y}}] \\ t_r = t - \frac{\sqrt{z^2 + a^2}}{c} \end{array} \right.$$

Problem 10.16 Show that the scalar potential of a point charge moving with constant velocity (Eq. 10.49) can be written more simply as

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{q}{R\sqrt{1 - v^2 \sin^2 \theta/c^2}}, \quad (10.51)$$

where $\mathbf{R} \equiv \mathbf{r} - \mathbf{v}t$ is the vector from the *present (!)* position of the particle to the field point \mathbf{r} , and θ is the angle between \mathbf{R} and \mathbf{v} (Fig. 10.9). Note that for nonrelativistic velocities ($v^2 \ll c^2$),

$$V(\mathbf{r}, t) \approx \frac{1}{4\pi\epsilon_0} \frac{q}{R}.$$

$$\begin{aligned} & (c^2 t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2) \\ &= [c^2 t - (\mathbf{R} + \mathbf{v}t) \cdot \mathbf{v}]^2 + (c^2 - v^2)[(\mathbf{R} + \mathbf{v}t) \cdot (\mathbf{R} + \mathbf{v}t) - c^2 t^2] \\ &= (c^2 t - Rv \cos \theta - v^2 t)^2 + (c^2 - v^2)(R^2 + 2Rvt \cos \theta + v^2 t^2 - c^2 t^2) \\ &= [(c^2 - v^2)t - Rv \cos \theta]^2 + (c^2 - v^2)(R^2 + 2Rvt \cos \theta) - [(c^2 - v^2)t]^2 \\ &= [(c^2 - v^2)t]^2 - 2(c^2 - v^2)tRv \cos \theta + (Rv \cos \theta)^2 \\ &\quad + (c^2 - v^2)(R^2 + 2Rvt \cos \theta) - [(c^2 - v^2)t]^2 \\ &= R^2 [(v \cos \theta)^2 + c^2 - v^2] \end{aligned}$$

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{(c^2 t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}}, \quad (10.49)$$

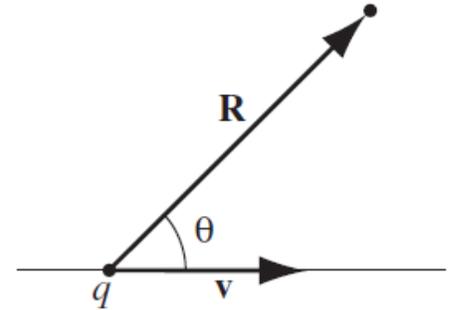


FIGURE 10.9

$$\begin{aligned} V(\mathbf{r}, t) &= \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{(c^2 t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2 t^2)}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{R^2 [(v \cos \theta)^2 + c^2 - v^2]}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R\sqrt{[1 + (v \cos \theta/c)^2 - (v/c)^2]}} \\ &= \frac{1}{4\pi\epsilon_0} \frac{q}{R\sqrt{[1 - (v \sin \theta/c)^2]}} \end{aligned}$$

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)}}, \quad (10.49)$$

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \frac{qc\mathbf{v}}{\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)}}. \quad (10.50)$$

$$\nabla \cdot \mathbf{A} = -\mu_0\epsilon_0 \frac{\partial V}{\partial t}. \quad (10.12)$$

Problem 10.27 Check that the potentials of a point charge moving at constant velocity (Eqs. 10.49 and 10.50) satisfy the Lorenz gauge condition (Eq. 10.12).

$$\begin{aligned} \nabla \cdot \mathbf{A} &= \frac{\mu_0}{4\pi} qc \left\{ \left[(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2) \right]^{-1/2} \nabla \cdot \mathbf{v} + \mathbf{v} \cdot \nabla \left[(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2) \right]^{-1/2} \right\} \\ &= \frac{\mu_0}{4\pi} qc\mathbf{v} \cdot \left\{ \left(-\frac{1}{2} \right) \left[(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2) \right]^{-3/2} \left[2(c^2t - \mathbf{r} \cdot \mathbf{v})\nabla(-\mathbf{r} \cdot \mathbf{v}) + (c^2 - v^2)\nabla(r^2) \right] \right\} \\ &= \frac{\mu_0}{8\pi} \frac{qc}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3} \mathbf{v} \cdot \left[2(c^2t - \mathbf{r} \cdot \mathbf{v})\nabla(\mathbf{r} \cdot \mathbf{v}) - (c^2 - v^2)\nabla(r^2) \right] \\ \nabla \cdot \mathbf{A} &= \frac{\mu_0}{8\pi} \frac{qc}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3} \mathbf{v} \cdot \left[2(c^2t - \mathbf{r} \cdot \mathbf{v})\mathbf{v} - (c^2 - v^2)2\mathbf{r} \right] \\ &= \frac{\mu_0 qc}{4\pi} \frac{(c^2t - \mathbf{r} \cdot \mathbf{v})v^2 - (c^2 - v^2)\mathbf{v} \cdot \mathbf{r}}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3} \\ &= \frac{\mu_0 qc^3}{4\pi} \frac{v^2t - \mathbf{v} \cdot \mathbf{r}}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3} \end{aligned}$$

$$\begin{aligned} \nabla(\mathbf{r} \cdot \mathbf{v}) &= \underbrace{(\mathbf{r} \cdot \nabla)\mathbf{v}}_{=0} + (\mathbf{v} \cdot \nabla)\mathbf{r} + \mathbf{r} \times \underbrace{(\nabla \times \mathbf{v})}_{=0} + \mathbf{v} \times \underbrace{(\nabla \times \mathbf{r})}_{=0} = \mathbf{v} \\ \nabla(r^2) &= \nabla(\mathbf{r} \cdot \mathbf{r}) = 2(\mathbf{r} \cdot \nabla)\mathbf{r} + 2\mathbf{r} \times \underbrace{(\nabla \times \mathbf{r})}_{=0} = 2\mathbf{r} \end{aligned}$$

Problem 10.27 Check that the potentials of a point charge moving at constant velocity (Eqs. 10.49 and 10.50) satisfy the Lorenz gauge condition (Eq. 10.12).

$$V(\mathbf{r}, t) = \frac{1}{4\pi\epsilon_0} \frac{qc}{\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)}}, \quad (10.49)$$

$$\mathbf{A}(\mathbf{r}, t) = \frac{\mu_0}{4\pi} \frac{qc\mathbf{v}}{\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)}}. \quad (10.50)$$

$$\boxed{\nabla \cdot \mathbf{A} = -\mu_0\epsilon_0 \frac{\partial V}{\partial t}}. \quad (10.12)$$

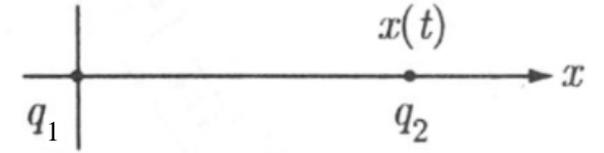
$$\begin{aligned} \frac{\partial V}{\partial t} &= \frac{qc}{4\pi\epsilon_0} \frac{\partial}{\partial t} \left[(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2) \right]^{-1/2} \\ &= \frac{qc}{4\pi\epsilon_0} \left(-\frac{1}{2} \right) \left[(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2) \right]^{-3/2} \left[2(c^2t - \mathbf{r} \cdot \mathbf{v}) \left(c^2 - \frac{\partial \mathbf{r} \cdot \mathbf{v}}{\partial t} \right) + (c^2 - v^2) \left(2\mathbf{r} \cdot \frac{\partial \mathbf{r}}{\partial t} - 2c^2t \right) \right] \\ &= -\frac{qc^3}{4\pi\epsilon_0} \frac{v^2t - \mathbf{r} \cdot \mathbf{v}}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3} \\ &\Rightarrow -\mu_0\epsilon_0 \frac{\partial V}{\partial t} = \nabla \cdot \mathbf{A} \end{aligned}$$

$$\nabla \cdot \mathbf{A} = \frac{\mu_0 qc^3}{4\pi} \frac{v^2t - \mathbf{v} \cdot \mathbf{r}}{\left(\sqrt{(c^2t - \mathbf{r} \cdot \mathbf{v})^2 + (c^2 - v^2)(r^2 - c^2t^2)} \right)^3}$$

Problem 10.28 One particle, of charge q_1 , is held at rest at the origin. Another particle, of charge q_2 , approaches along the x axis, in hyperbolic motion:

$$x(t) = \sqrt{b^2 + (ct)^2};$$

it reaches the closest point, b , at time $t = 0$, and then returns out to infinity.



(a) What is the force F_2 on q_2 (due to q_1) at time t ?

$$\mathbf{F}_2 = \frac{q_1 q_2}{4\pi\epsilon_0} \frac{1}{b^2 + c^2 t^2} \hat{\mathbf{x}}$$

(b) What total impulse ($I_2 = \int_{-\infty}^{\infty} F_2 dt$) is delivered to q_2 by q_1 ?

$$I_2 = \int_{-\infty}^{\infty} F_2 dt = \int_{-\infty}^{\infty} \frac{q_1 q_2}{4\pi\epsilon_0} \frac{1}{b^2 + c^2 t^2} dt = \frac{q_1 q_2}{4\pi\epsilon_0} \left[\frac{1}{bc} \tan^{-1} \frac{ct}{b} \right]_{-\infty}^{\infty} = \frac{q_1 q_2}{4\pi\epsilon_0} \frac{1}{bc} \left[\frac{\pi}{2} - \left(-\frac{\pi}{2} \right) \right] = \frac{q_1 q_2}{4\pi\epsilon_0} \frac{\pi}{bc}$$

Problem 10.28

(c) What is the force F_1 on q_1 (due to q_2) at time t ?

$$c(t - t_r) = |\mathbf{r} - \mathbf{w}(t_r)| = x(t_r) = \sqrt{b^2 + c^2 t_r^2} \Rightarrow t_r = \frac{c^2 t^2 - b^2}{2c^2 t} \quad (\text{Do not assume } v \text{ is constant!})$$

What are the fields on the axis to the left of the charge?

$$\begin{cases} x(t) = \sqrt{b^2 + c^2 t^2} \\ v(t) = \frac{c^2 t}{\sqrt{b^2 + c^2 t^2}} = \frac{c^2 t}{x} \end{cases} \Rightarrow \begin{cases} x(t_r) = c(t - t_r) = \frac{c^2 t^2 + b^2}{2ct} \\ v(t_r) = \frac{c^2 t_r}{x(t_r)} = \frac{c^2 t^2 - b^2}{2t} \frac{2ct}{c^2 t^2 + b^2} = c \left(\frac{c^2 t^2 - b^2}{c^2 t^2 + b^2} \right) \end{cases}$$

$$\mathbf{E} = \begin{cases} \frac{q_2}{4\pi\epsilon_0} \frac{1}{r^2} \left(\frac{c+v}{c-v} \right) \hat{\mathbf{x}} & \text{, to the right of the charge} \\ -\frac{q_2}{4\pi\epsilon_0} \frac{1}{r^2} \left(\frac{c-v}{c+v} \right) \hat{\mathbf{x}} & \text{, to the left of the charge} \end{cases}$$

$$\mathbf{E} = -\frac{q_2}{4\pi\epsilon_0} \frac{1}{x(t_r)^2} \left(\frac{c-v}{c+v} \right) \hat{\mathbf{x}} = -\frac{q_2}{4\pi\epsilon_0} \frac{1}{x(t_r)^2} \left(1 - \frac{c^2 t^2 - b^2}{c^2 t^2 + b^2} \right) \left(1 + \frac{c^2 t^2 - b^2}{c^2 t^2 + b^2} \right)^{-1} \hat{\mathbf{x}}$$

$$= -\frac{q_2}{4\pi\epsilon_0} \left(\frac{c^2 t^2 + b^2}{2ct} \right)^{-2} \frac{2b^2}{c^2 t^2 + b^2} \frac{c^2 t^2 + b^2}{2c^2 t^2} \hat{\mathbf{x}}$$

$$= -\frac{q_2}{4\pi\epsilon_0} \frac{4c^2 t^2}{(c^2 t^2 + b^2)^2} \frac{b^2}{c^2 t^2} \hat{\mathbf{x}}$$

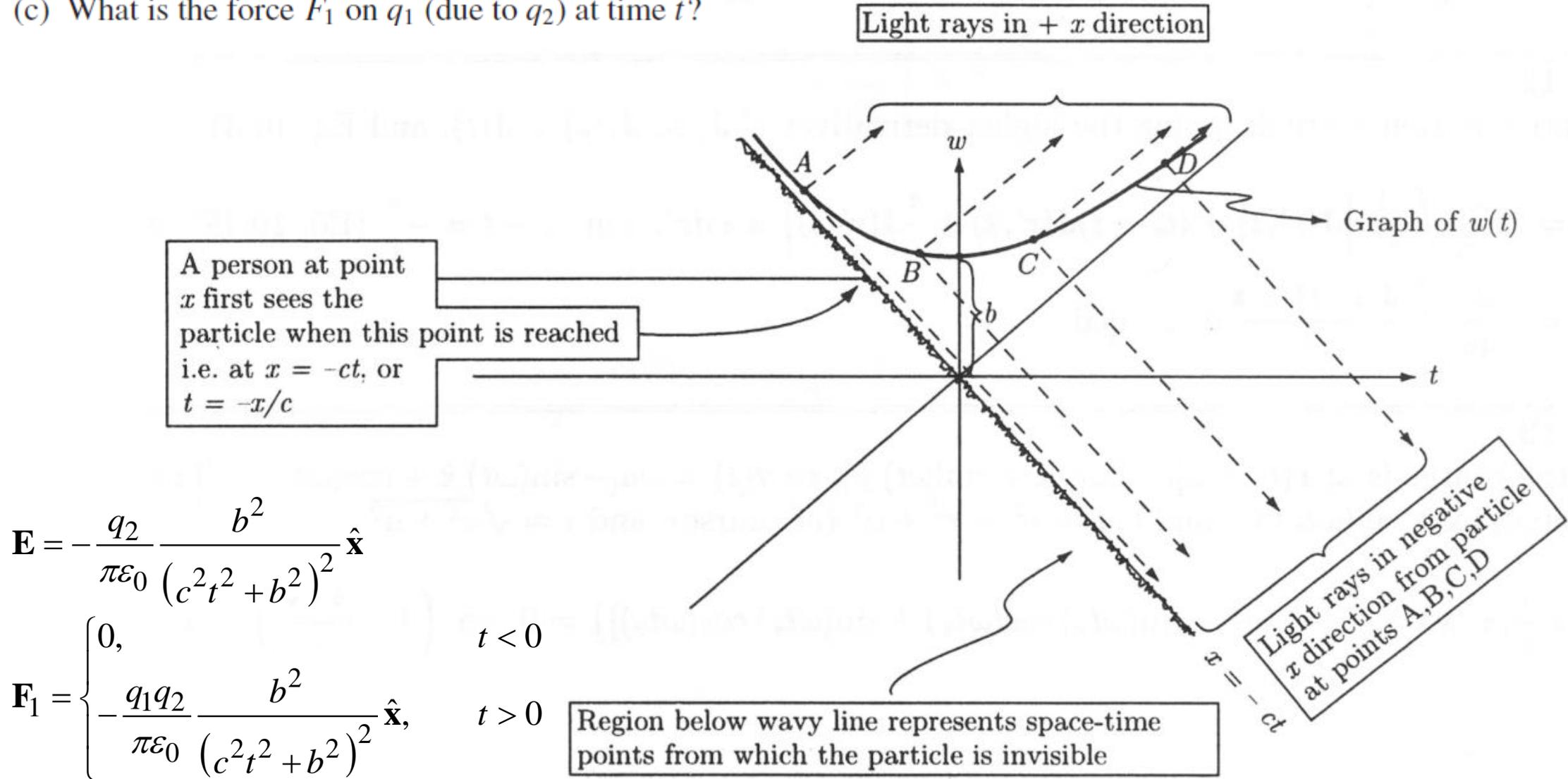
$$= -\frac{q_2}{\pi\epsilon_0} \frac{b^2}{(c^2 t^2 + b^2)^2} \hat{\mathbf{x}}$$

Problem 10.20 Suppose a point charge q is constrained to move along the x axis. Show that the fields at points on the axis to the right of the charge are given by

$$\mathbf{E} = \frac{q}{4\pi\epsilon_0} \frac{1}{r^2} \left(\frac{c+v}{c-v} \right) \hat{\mathbf{x}}, \quad \mathbf{B} = \mathbf{0}.$$

Problem 10.28

(c) What is the force F_1 on q_1 (due to q_2) at time t ?



$$\mathbf{E} = -\frac{q_2}{\pi\epsilon_0} \frac{b^2}{(c^2t^2 + b^2)^2} \hat{\mathbf{x}}$$

$$\mathbf{F}_1 = \begin{cases} 0, & t < 0 \\ -\frac{q_1q_2}{\pi\epsilon_0} \frac{b^2}{(c^2t^2 + b^2)^2} \hat{\mathbf{x}}, & t > 0 \end{cases}$$

Problem 10.28

(d) What total impulse ($I_1 = \int_{-\infty}^{\infty} F_1 dt$) is delivered to q_1 by q_2 ? [Hint: It might help to review Prob. 10.17 before doing this integral. Answer: $I_2 = -I_1 = q_1 q_2 / 4\epsilon_0 bc$]

$$I_2 = \frac{q_1 q_2}{4\pi\epsilon_0} \frac{\pi}{bc}$$

$$I_1 = -\frac{q_1 q_2}{\pi\epsilon_0} b^2 \int_0^{\infty} \frac{1}{(c^2 t^2 + b^2)^2} dt = -\frac{q_1 q_2}{\pi\epsilon_0} b^2 \frac{1}{2b^2} \left(\frac{t}{b^2 + c^2 t^2} \Big|_0^{\infty} + \int_0^{\infty} \frac{1}{b^2 + c^2 t^2} dt \right) = -\frac{q_1 q_2}{2\pi\epsilon_0} \left(\frac{1}{bc} \frac{\pi}{2} \right) = -\frac{q_1 q_2}{4\pi\epsilon_0} \frac{\pi}{bc}$$

$$\Rightarrow I_2 = -I_1 = \frac{q_1 q_2}{4\pi\epsilon_0} \frac{\pi}{bc}$$