

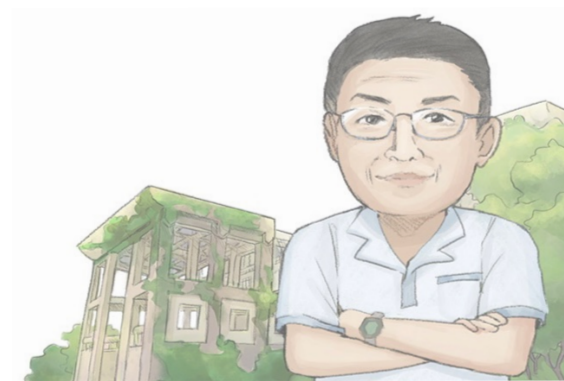


Quiz 2

Fall, 2023

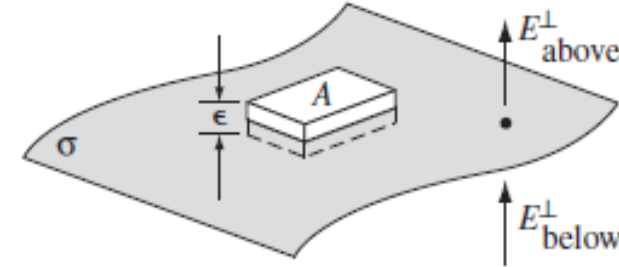
Tien-Fu Yang

Department of Physics, National Tsing Hua University



Electrostatic Boundary Conditions: Normal

The electric field is not continuous at a surface with charge density σ . Why?

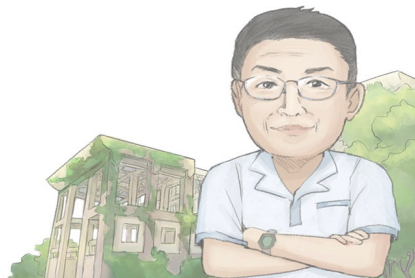


Consider a Gaussian pillbox.

Gauss's law states that
$$\oint_S \mathbf{E} \cdot d\mathbf{a} = \frac{Q_{\text{enc}}}{\epsilon_0} = \frac{\sigma A}{\epsilon_0}$$

The sides of the pillbox contribute nothing to the flux, in the limit as the thickness ϵ goes to zero.

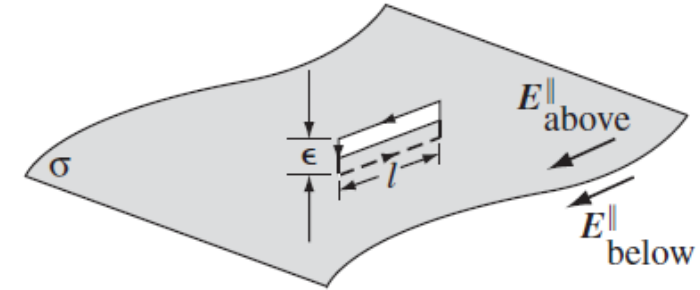
$$(E_{\text{above}}^{\perp} - E_{\text{below}}^{\perp})A = \frac{\sigma A}{\epsilon_0} \Rightarrow (E_{\text{above}}^{\perp} - E_{\text{below}}^{\perp}) = \frac{\sigma}{\epsilon_0}$$



Electrostatic Boundary Conditions: **Tangential**

The tangential component of \mathbf{E} , by contract, is always continuous.

Consider a thin rectangular loop.

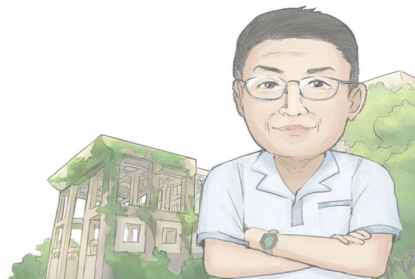


The curl of the electric field states that $\oint_P \mathbf{E} \cdot d\vec{\ell} = 0$

The ends give nothing (as $\epsilon \rightarrow 0$), and the sides give

$$(E_{\text{above}}^{\parallel} - E_{\text{below}}^{\parallel})\ell = 0 \quad \Rightarrow \quad E_{\text{above}}^{\parallel} = E_{\text{below}}^{\parallel}$$

$$\text{In short, } \mathbf{E}_{\text{above}} - \mathbf{E}_{\text{below}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$



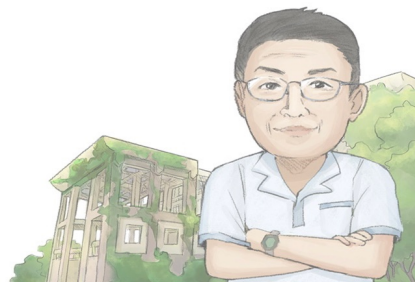
Boundary Conditions in terms of potential

$$\mathbf{E}_{\text{above}} - \mathbf{E}_{\text{below}} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}} \Rightarrow (\nabla V_{\text{above}} - \nabla V_{\text{below}}) = -\frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$$

$$\text{or } \left(\frac{\partial V_{\text{above}}}{\partial n} - \frac{\partial V_{\text{below}}}{\partial n} \right) = -\frac{\sigma}{\epsilon_0}$$

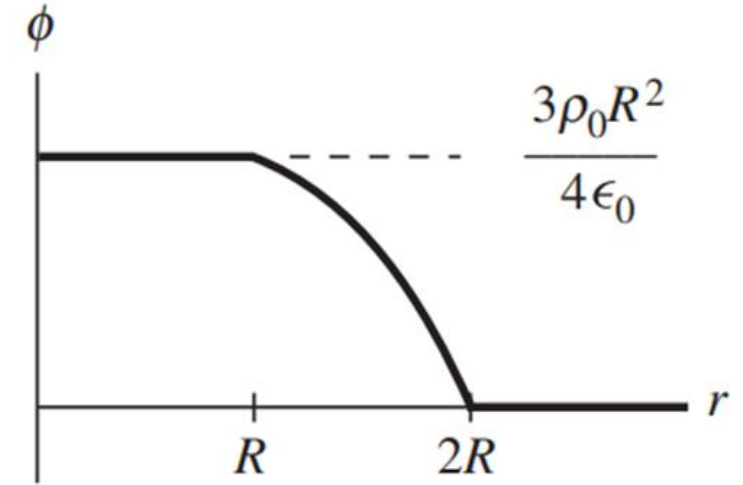
$$\text{where } \frac{\partial V_{\text{above}}}{\partial n} = \nabla V \cdot \hat{\mathbf{n}}$$

denotes the normal derivative of V .



A distribution of charge has cylindrical symmetry. As a function of the distance r from the symmetry axis, the electric potential is

$$\phi(r) = \begin{cases} \frac{3\rho_0 R^2}{4\epsilon_0} & (\text{for } r \leq R) \\ \frac{\rho_0}{4\epsilon_0} (4R^2 - r^2) & (\text{for } R < r < 2R) \\ 0 & (\text{for } r \geq 2R) \end{cases},$$

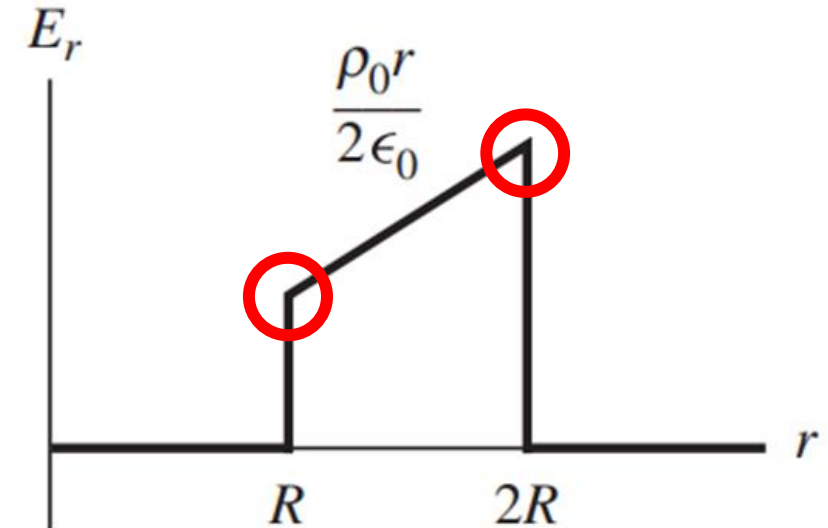


where ρ_0 is a quantity with the dimensions of volume charge density.

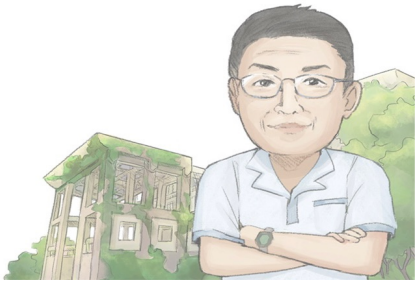
- Find and make rough plots of the electric field, for all values of r .
- Determine the charge distribution and explain the reasons for the discontinuities in the electric field.

$$\mathbf{E} = \begin{cases} -\nabla\phi = -\hat{\mathbf{r}} \frac{\partial}{\partial r} \left[\frac{\rho_0}{4\epsilon_0} (4R^2 - r^2) \right] = \frac{\rho_0 r}{2\epsilon_0} \hat{\mathbf{r}} & \text{for } R < r < 2R \\ 0, & \text{otherwise} \end{cases}$$

$$\Rightarrow \rho = \epsilon_0 \nabla \cdot \mathbf{E} = \epsilon_0 \frac{1}{r} \frac{\partial}{\partial r} \left[r \cdot \frac{\rho_0 r}{2\epsilon_0} \right] = \rho_0 \quad \begin{cases} \sigma_R = \frac{\rho_0 R}{2} \\ \sigma_{2R} = -\rho_0 R \end{cases}$$



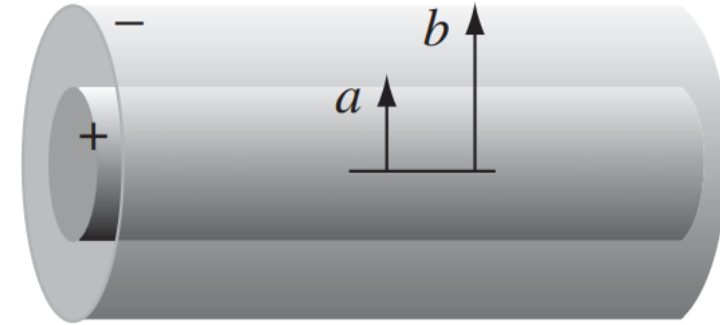
Discontinuities → Surface charges



From your charge distribution, calculate the total charge per unit length along the cylinder. Explain the result.

The cross-sectional area of the $R < r < 2R$ region is $\pi(2R)^2 - \pi R^2 = 3\pi R^2$. So the volume density ρ_0 in this region yields a charge in a length ℓ of the cylinder equal to

$$\rho_0 \ell (3\pi R^2) = 3\pi R^2 \rho_0 \ell.$$



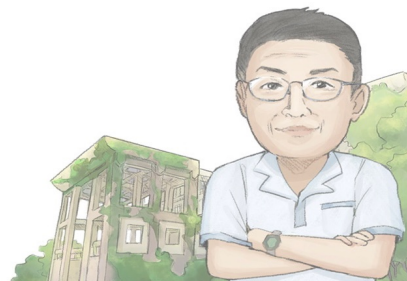
The surface at $r = R$ yields a charge in length ℓ equal to

$$\sigma_R (2\pi R) \ell = (\rho_0 R/2) (2\pi R) \ell = \pi R^2 \rho_0 \ell.$$

And the surface at $r = 2R$ yields a charge in length ℓ equal to

$$\sigma_{2R} (2\pi \cdot 2R) \ell = (-\rho_0 R) (4\pi R) \ell = -4\pi R^2 \rho_0 \ell.$$

Adding up the above three charges and dividing by , we see that the total charge per unit length is zero.





Homework Exercises (Chap.3)

Griffiths:

11, 13, 16, 20, 27, 43, 54, 56,

7, 8, 12, 19, 23, 28, 29, 32, 44, 49

